

Bielefeld University
Department of Physics

A. Weddemann, A. Auge,
F. Wittbracht, S. Herth, A. Hütten

Interactions of magnetic particles in a rotational magnetic field

D2 PHYSICS

Bielefeld University

Outline

2

A. Weddemann, A. Auge,
F. Wittbracht, S. Herth, A. Hütten
06.11.2008

1. Motivation
2. Governing equation
 1. Particle motion
 2. Technical realization - ALE-approach
 1. Second domain triangulation
3. Interactions of beads in fluids
 1. Simple system
 2. Comparison between magnetic and hydrodynamic forces
4. Conclusions and Outlook

Motivation

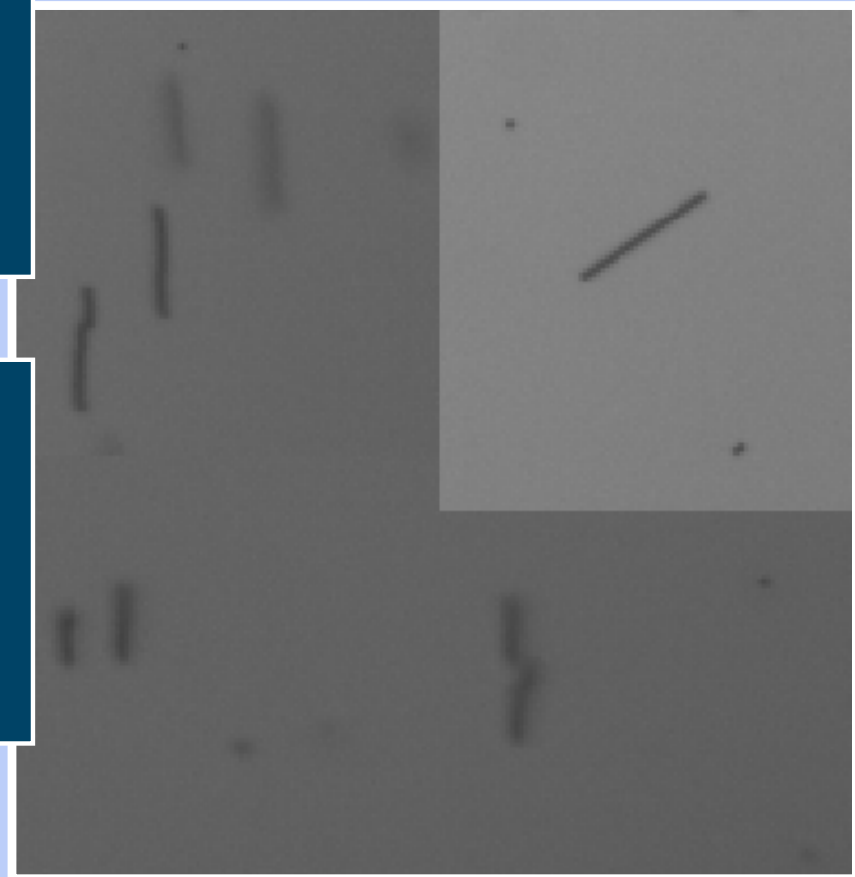
3

Experimental observations

Magnetic micro- or nanoparticles can interact very strongly:

Under the influence of an external homogenous magnetic field particles create chains

Question: Can magnetic interactions be neglected when modeling particles in microfluidic systems?

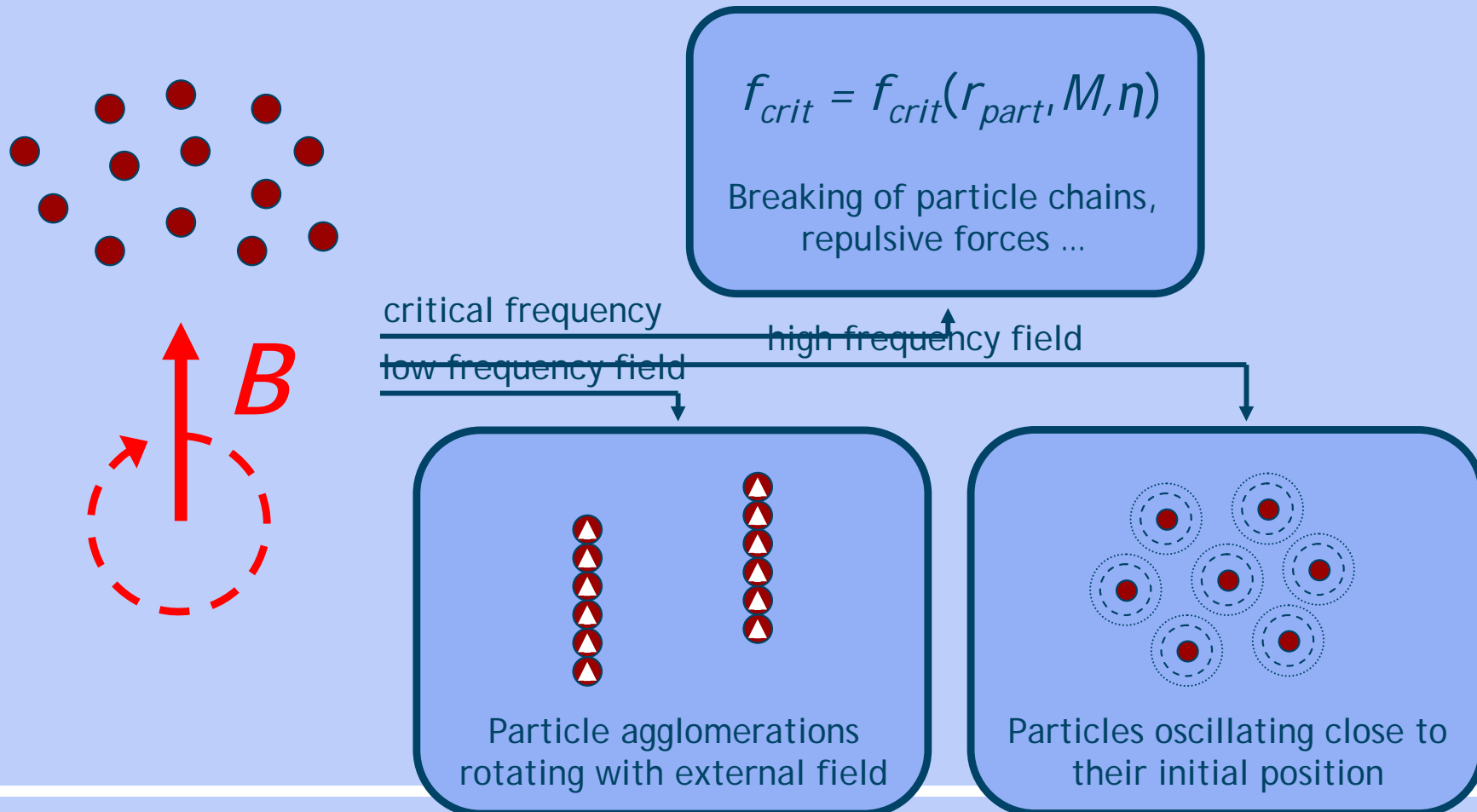


Motivation

4

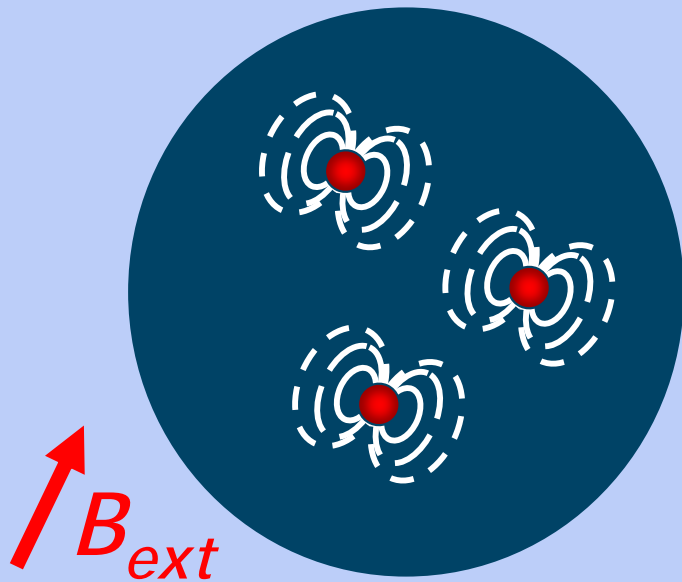
Model system

Expect ferro- or superparamagnetic particles in a solution, being manipulated by an external magnetic field.



Particle motion

Governing equations:



Particle motion:

$$M \frac{d}{dt} U(t) = F_{mag} + F_{visc} + F_{pen}$$

$$F_{mag} = \int_{particle} f dx = - \int_{particle} \text{grad} \langle M, B \rangle dx$$

F_{visc} □ viscous force term

F_{pen} □ force term preventing particles from overlapping

$U(t) = (v_x^{part_1}, v_y^{part_1}, v_x^{part_2}, \dots)^T$ □ velocity vector

$$M \square B_{ext} \quad |M| = M_s$$

$$\int_{\Omega_t} \langle \text{grad} \psi_{A_z}, \text{grad} A_z \rangle dx - \mu_0 \int_{\Omega_t} \left(M_y \frac{\partial \psi_{A_z}}{\partial x} - M_x \frac{\partial \psi_{A_z}}{\partial y} \right) dx = 0$$

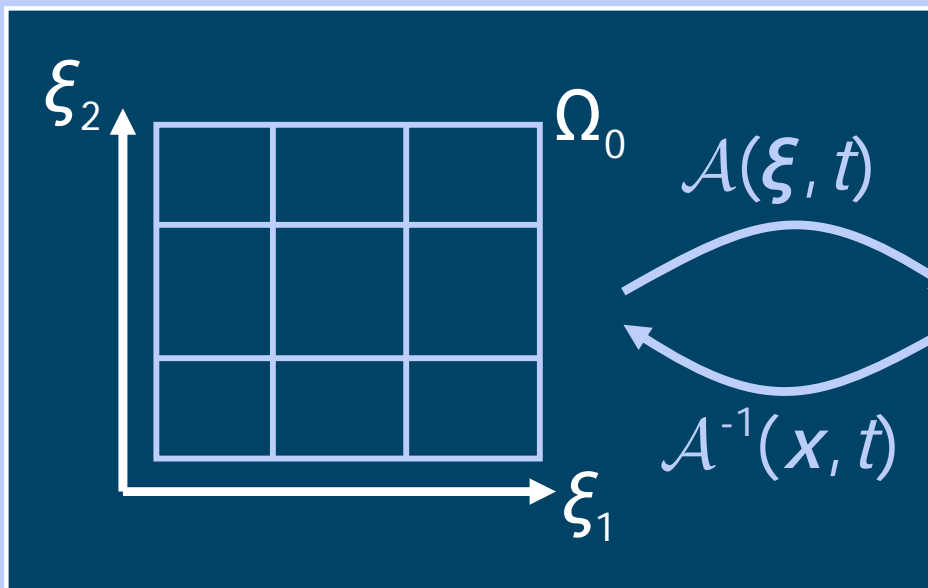
$$M \frac{d}{dt} U(t) = F_{mag} + F_{visc} + F_{pen}$$

Particle movement requires mesh displacement

→ ALE-formalism

ALE-formulation

The basic idea of ALE-methods is to use different coordinate systems, a reference and a spatial system.



Calculation transformed to reference system

Example:

$$\frac{\partial u}{\partial t}(x, t) + \mathcal{L}[u](x, t) = 0$$

----- · weak formulation -----

$$\int_{\Omega_t} \psi(x, t) \cdot \frac{\partial u}{\partial t}(x, t) dx + \int_{\Omega_t} \psi(x, t) \cdot \mathcal{L}[u](x, t) dx = 0$$

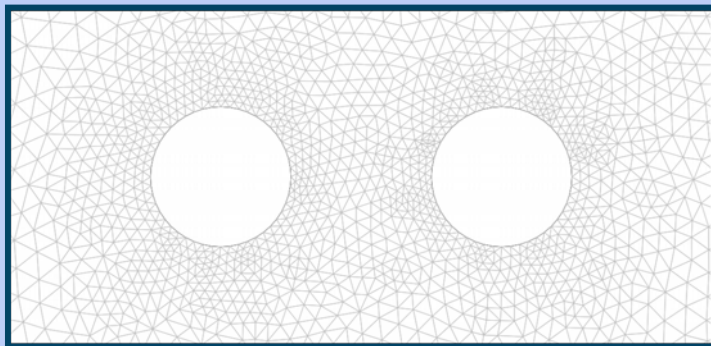
----- · domain transformation -----

$$\int_{\Omega_0} \psi(\mathcal{A}(\xi, t)) \cdot \frac{\partial u}{\partial t}(\mathcal{A}(\xi, t), t) \cdot \det J_{\mathcal{A}_t}(\xi, t) d\xi + \int_{\Omega_0} \psi(\mathcal{A}(\xi, t)) \cdot \mathcal{L}[u](\mathcal{A}(\xi, t), t) \cdot \det J_{\mathcal{A}_t}(\xi, t) d\xi = 0$$

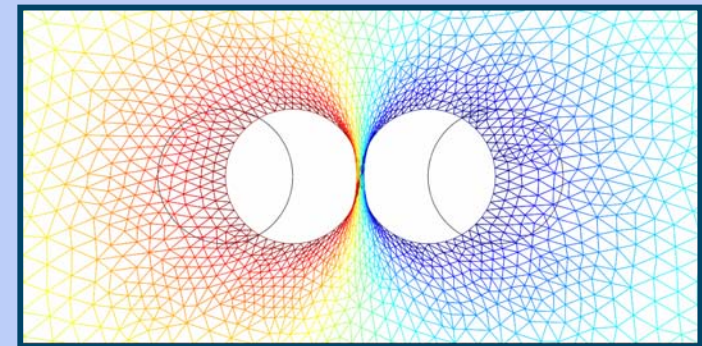
ALE-formulation

Limitations of ALE-methods:

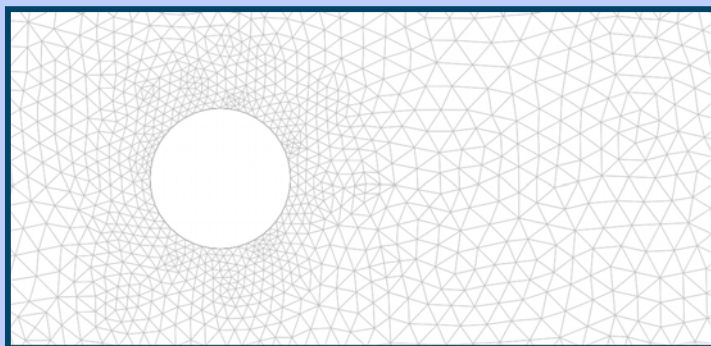
- Topological changes:



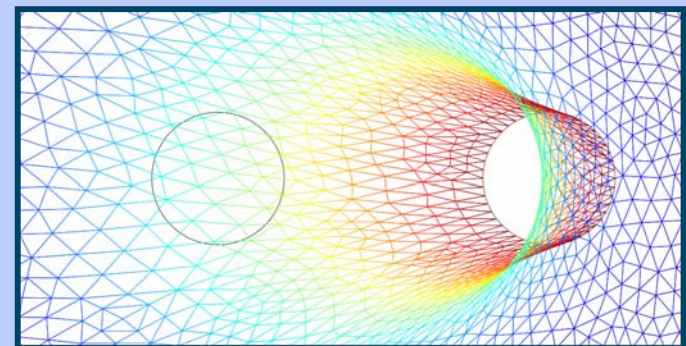
particles moving towards each other



- Very strong displacements:



particle moving too far in one direction

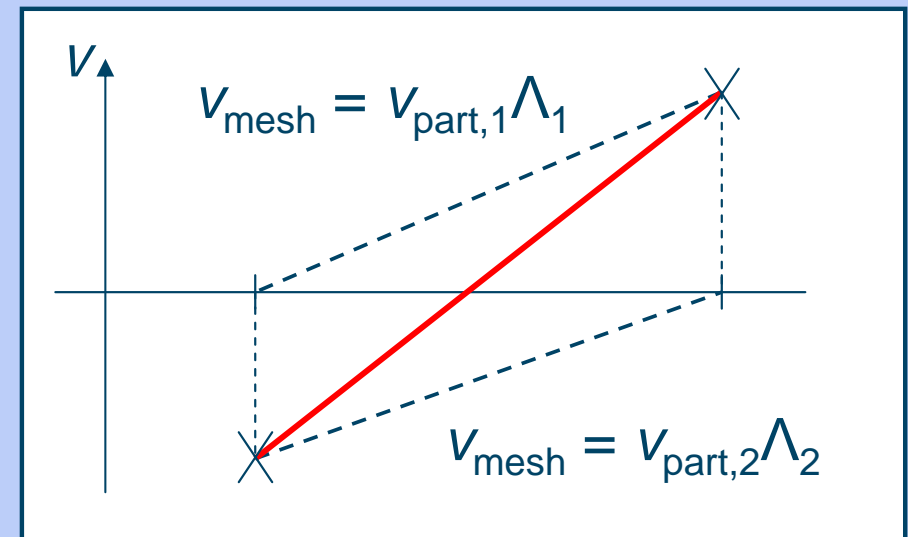
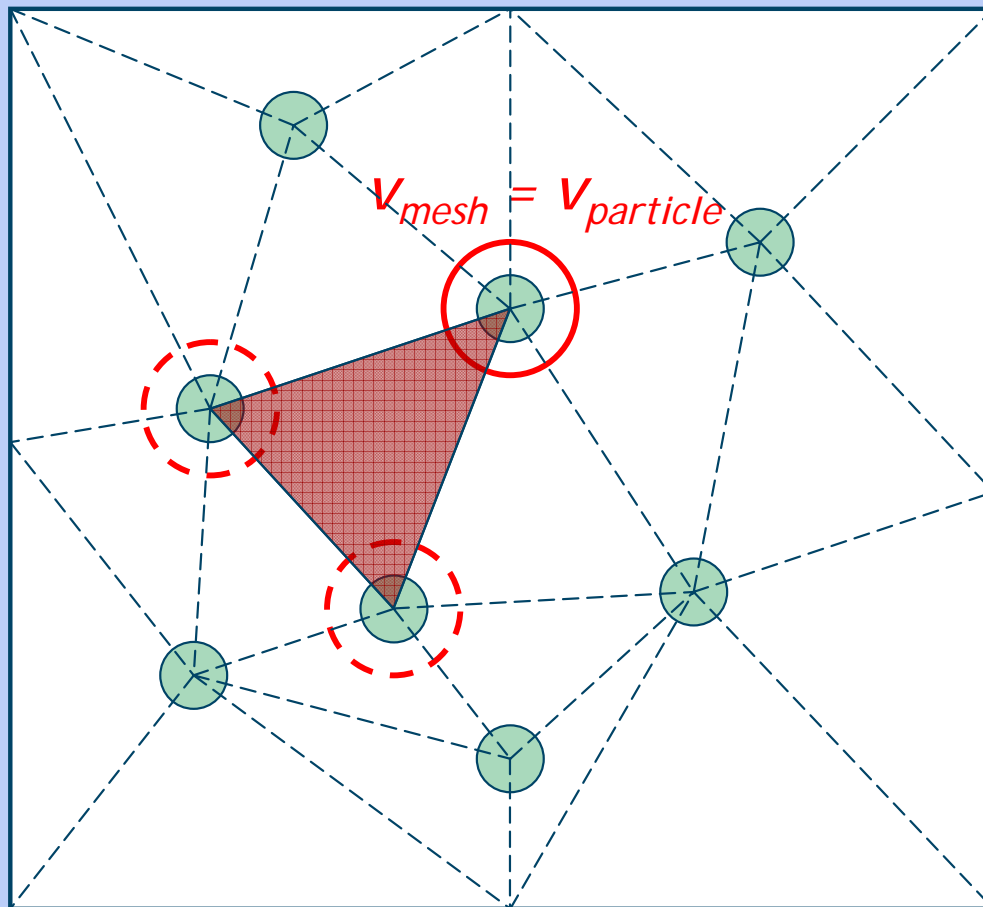


ALE-formulation

8

Governing equations

Calculation of ALE-mesh-displacement:

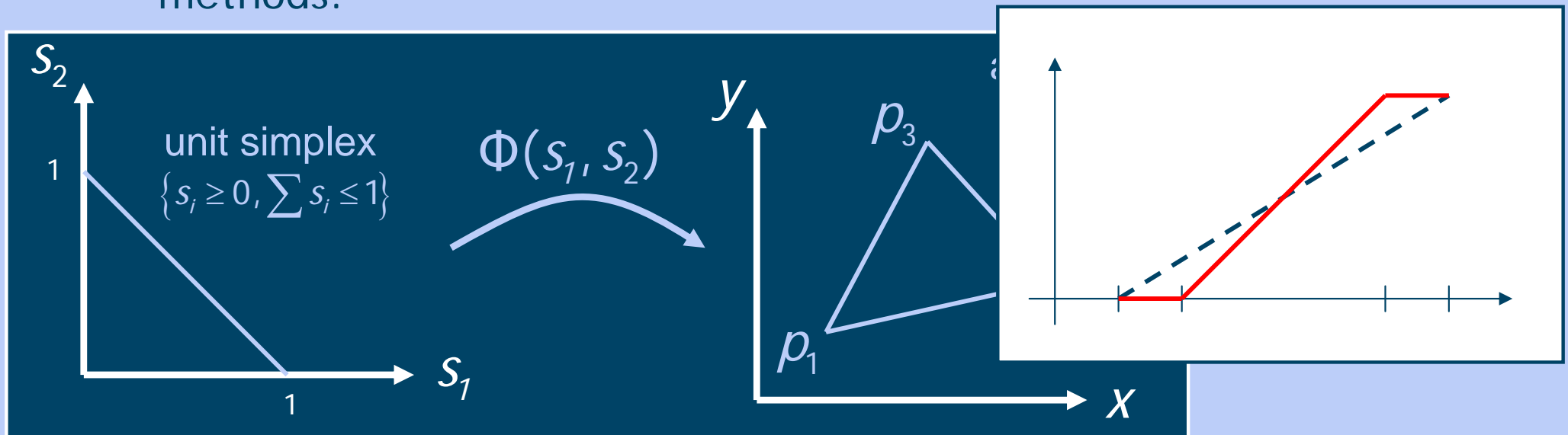


$$v_{mesh} = \sum_{\text{nodes}} v_{part,i} \Lambda_i$$

second FEM-triangulation
with linear basis set Λ

Second domain triangulation

The parameter functions Λ can be calculated by standard FEM-methods:



with affine linear mapping

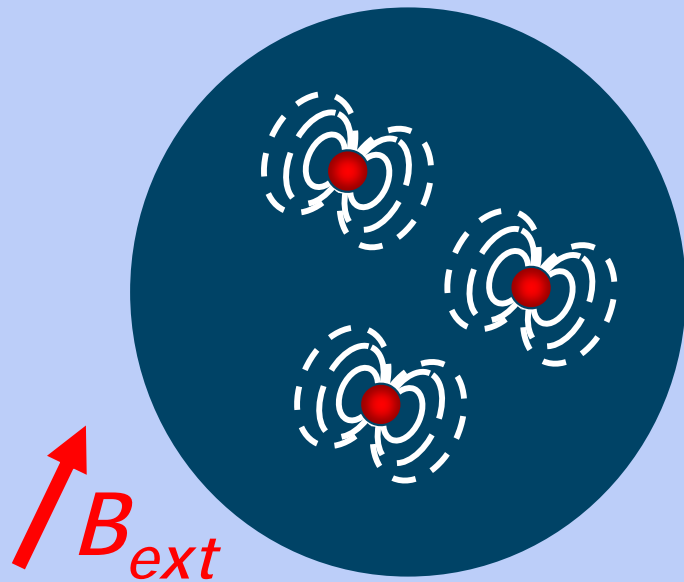
$$\Phi_{x_1 x_2 x_3}(s_1, s_2) = \begin{pmatrix} x_1 \\ y_1 \end{pmatrix} + s_1 \begin{pmatrix} x_2 - x_1 \\ y_2 - y_1 \end{pmatrix} + s_2 \begin{pmatrix} x_3 - x_1 \\ y_3 - y_1 \end{pmatrix}$$

$$\Lambda(x) = \Lambda(\Phi_{x_1 x_2 x_3}^{-1}(x)) = \frac{(x_3 - x)(y_3 - y_2) - (x_3 - x_2)(y_3 - y)}{(x_2 - x_1)(y_3 - y_1) - (x_3 - x_1)(y_2 - y_1)}$$

$$f(\Lambda, \theta_1, \theta_2) = \frac{\Lambda - \theta_1}{1 - (\theta_1 + \theta_2)} \cdot \Theta(\Lambda - \theta_1) \cdot \Theta(1 - \Lambda - \theta_2) + \Theta(\Lambda - (1 - \theta_2))$$

Particle motion

Governing equations:



Particles induce fluid flow:

$$\rho \frac{\partial u}{\partial t} + \rho(u \nabla) u = -\text{grad } p + \eta \Delta u + \rho f$$

$$\text{div } u = 0$$

Total mesh displacement:

$$\Delta r = \sum_i (r_i - \xi_i) \cdot f(\Lambda_i(r), \theta_1, \theta_2)$$

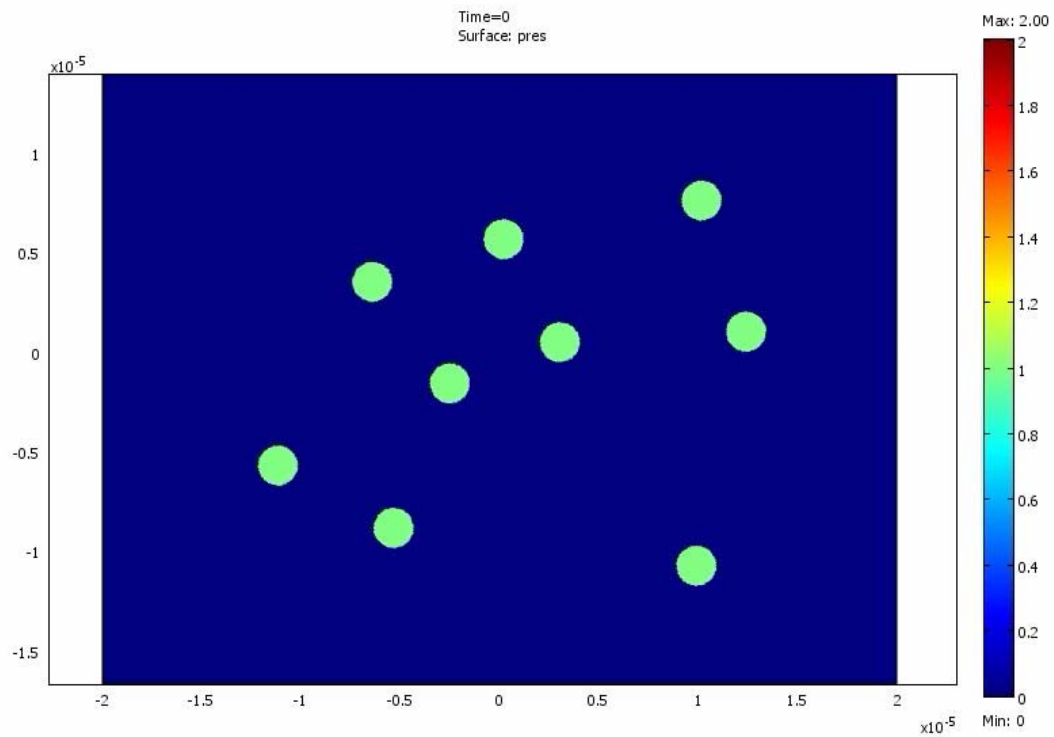
Additional remeshing condition:

$$\min_{T \in \mathcal{T}} \text{qual } T < \sigma$$

$$M \square B_{ext} \quad |M| = M_s$$

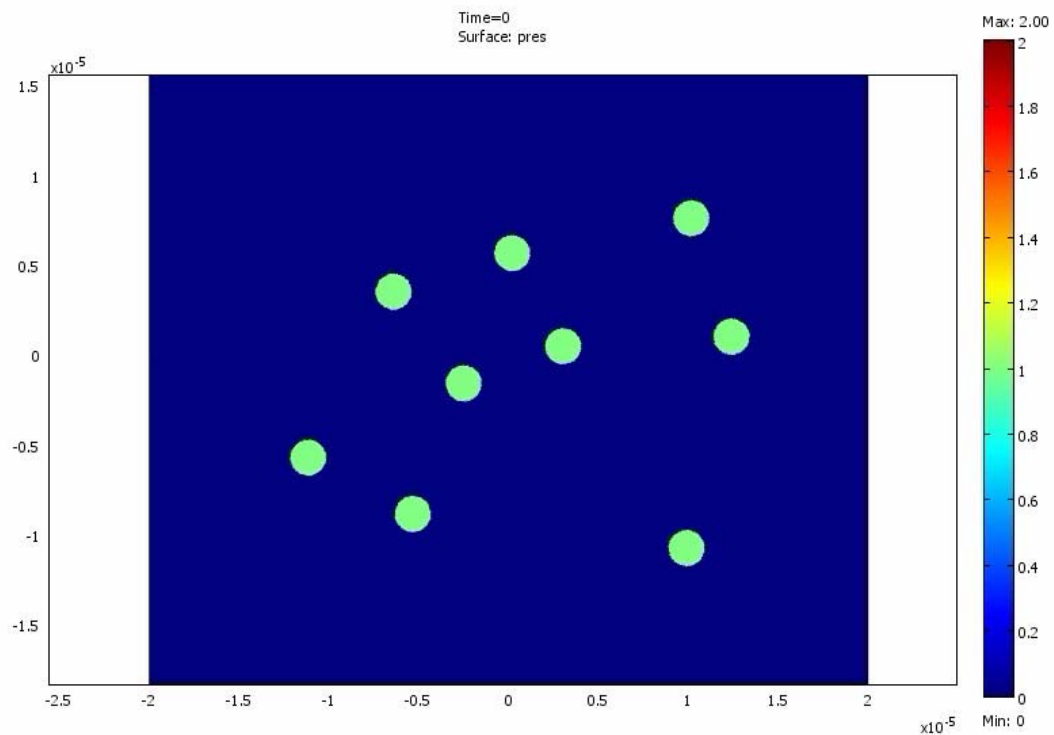
$$\int_{\Omega_t} \langle \text{grad } \psi_{A_z}, \text{grad } A_z \rangle dx - \mu_0 \int_{\Omega_t} \left(M_y \frac{\partial \psi_{A_z}}{\partial x} - M_x \frac{\partial \psi_{A_z}}{\partial y} \right) dx = 0$$

$$M \frac{d}{dt} U(t) = F_{mag} + F_{visc} + F_{pen}$$



Interactions of beads in fluids

$$f = f_0$$



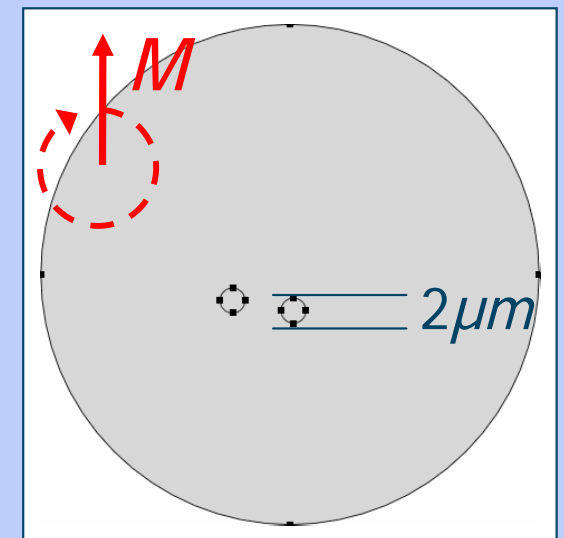
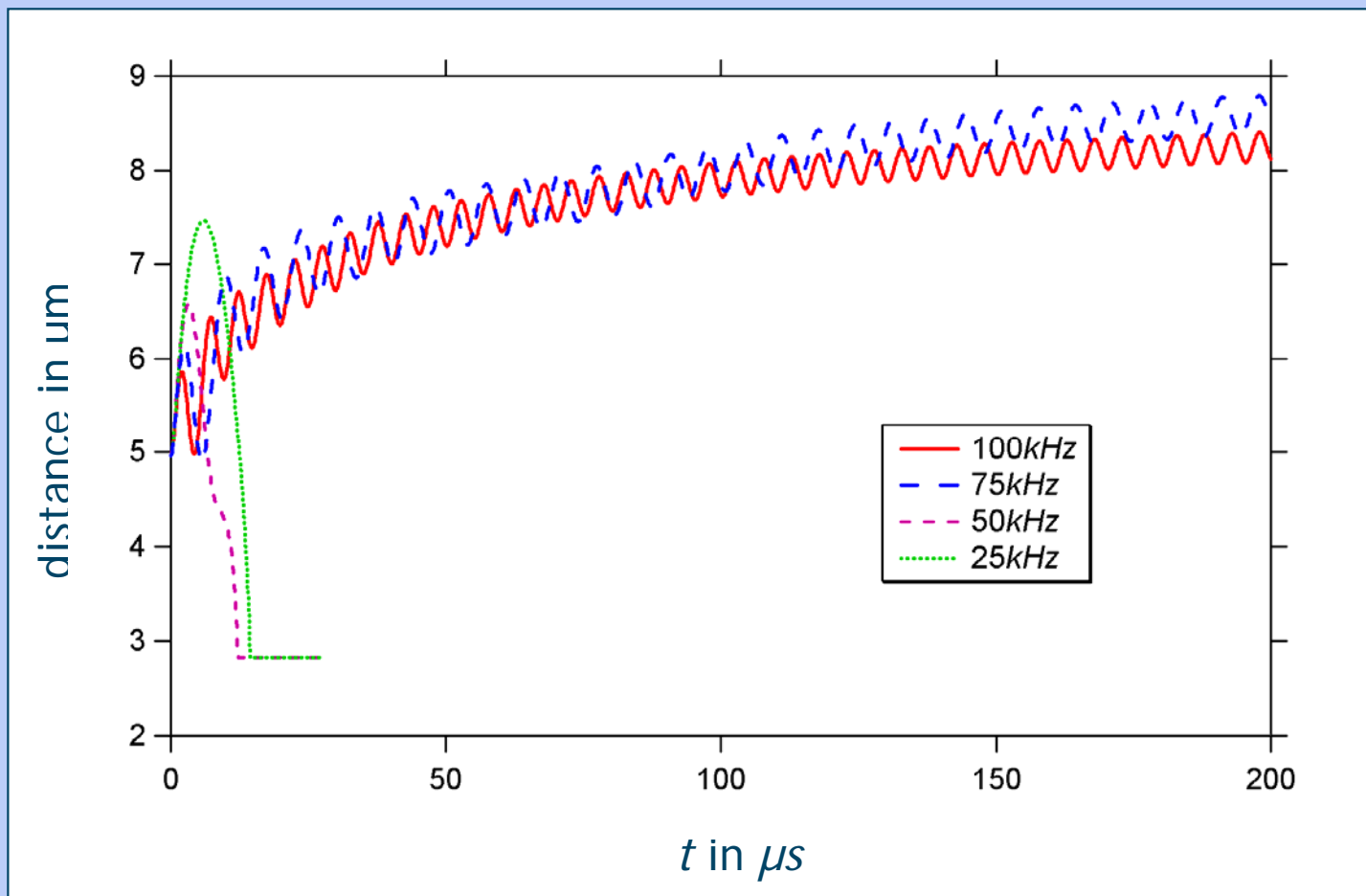
$$f = 1.5f_0$$

Different phenomena:

- chain creation
- particles oscillating against each other

Interactions of beads in fluids

Frequency dependence for different initial conditions:



$$M_s = 1000 \text{ kAm}^{-1}$$

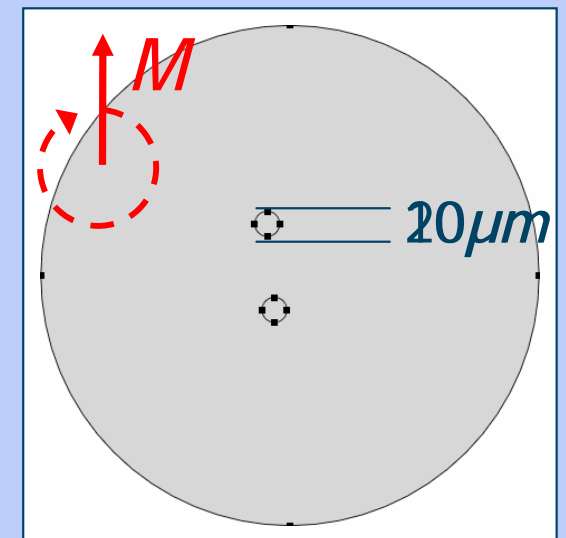
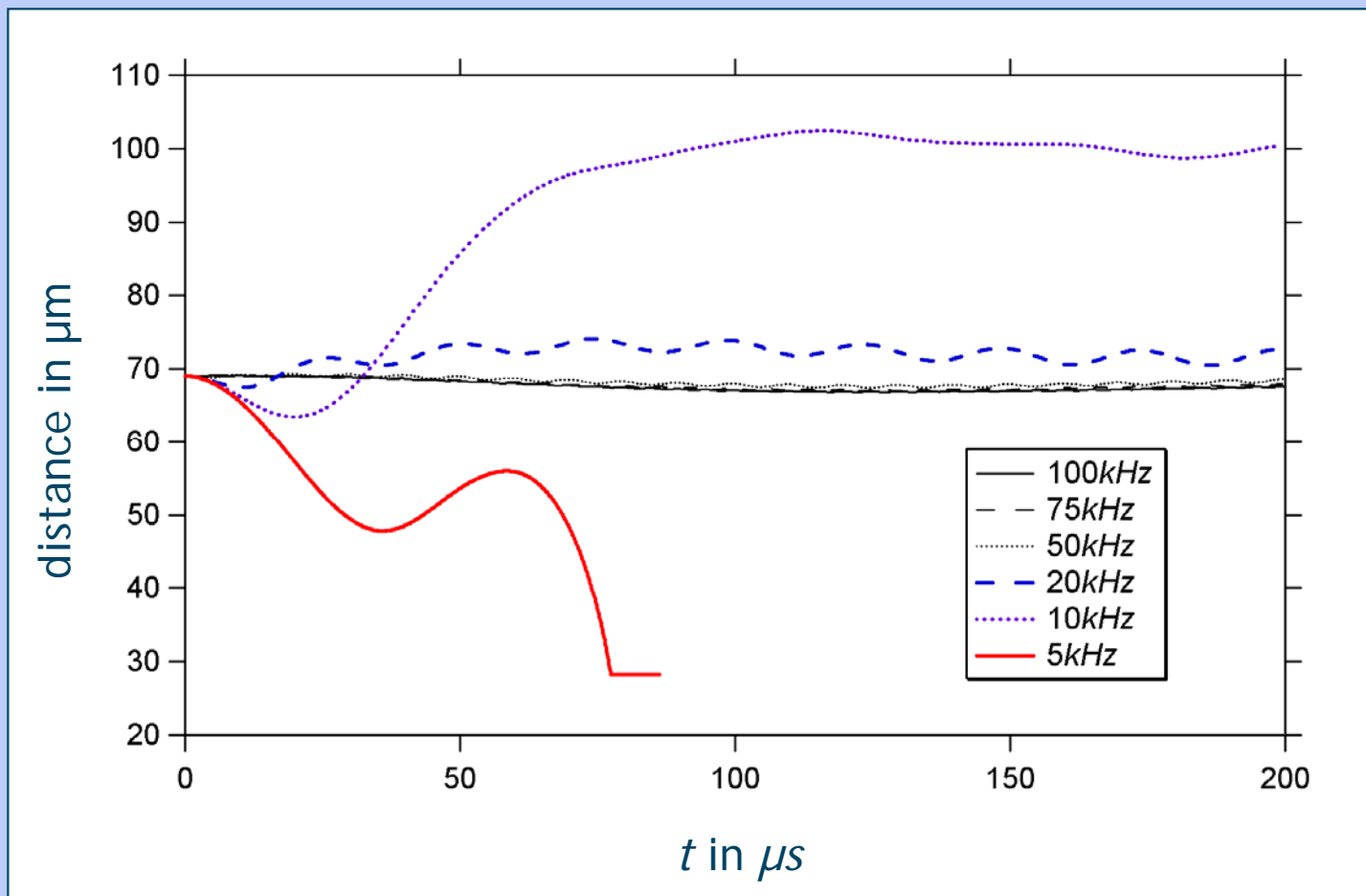
$$r = 1 \mu\text{m}$$

Interactions of beads in fluids

13

Interactions of beads in fluids

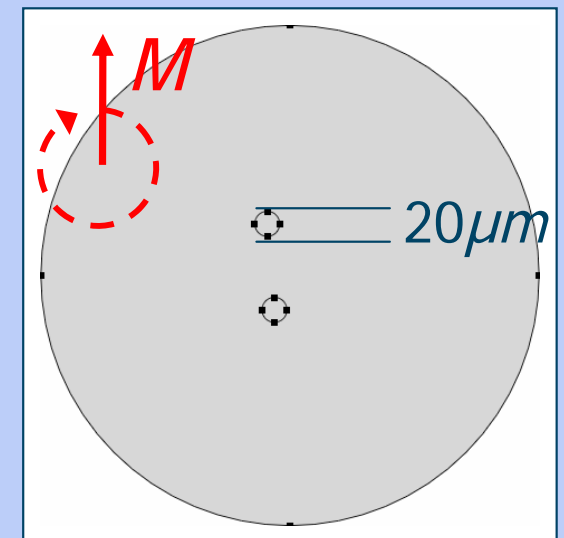
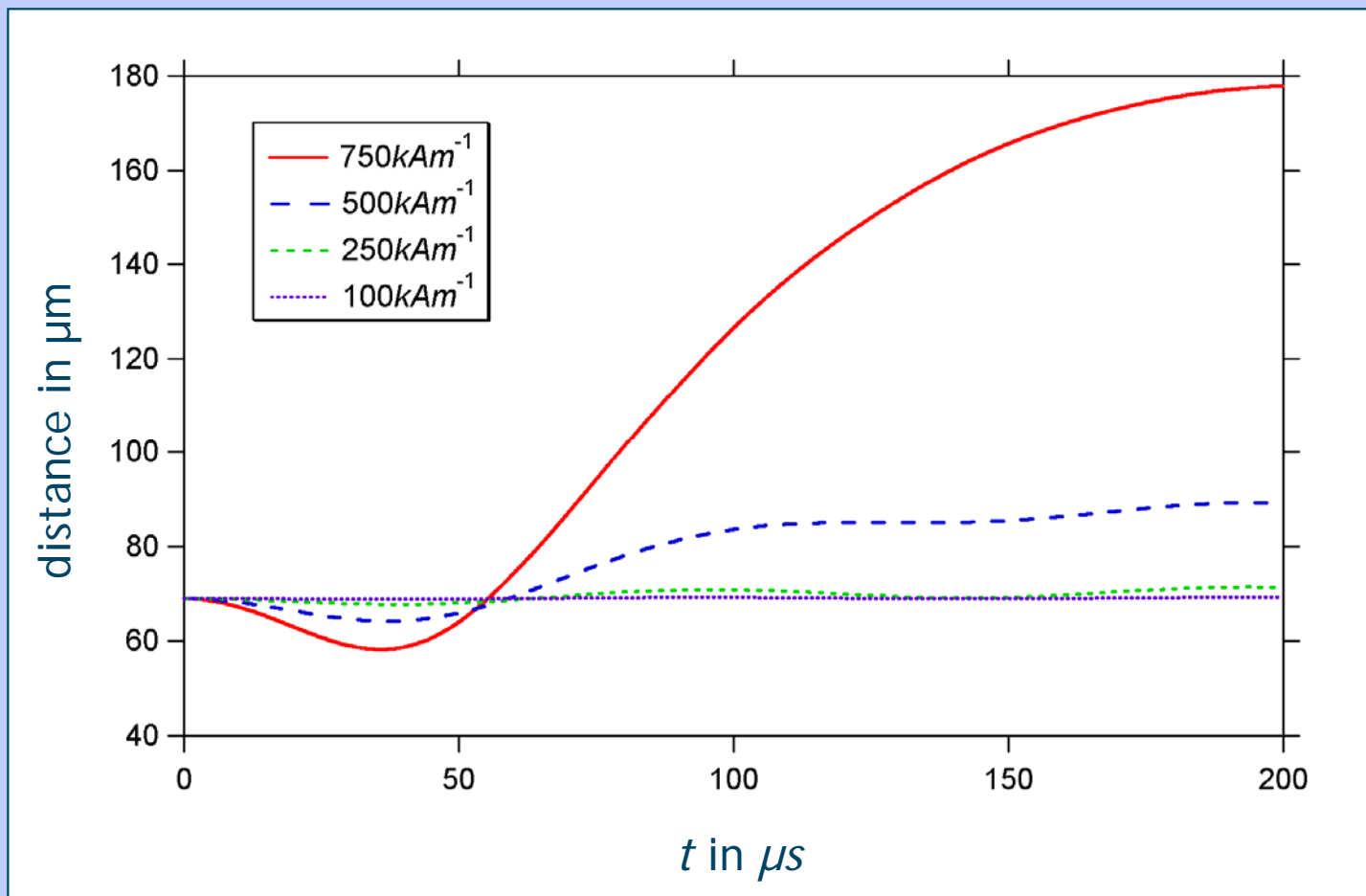
Frequency dependence for different particle diameters:



$$M_s = 1000 \text{ kAm}^{-1}$$
$$r = 20 \mu\text{m}$$

Interactions of beads in fluids

Frequency dependence for different particle magnetizations:



$$f = 20\text{kHz}$$

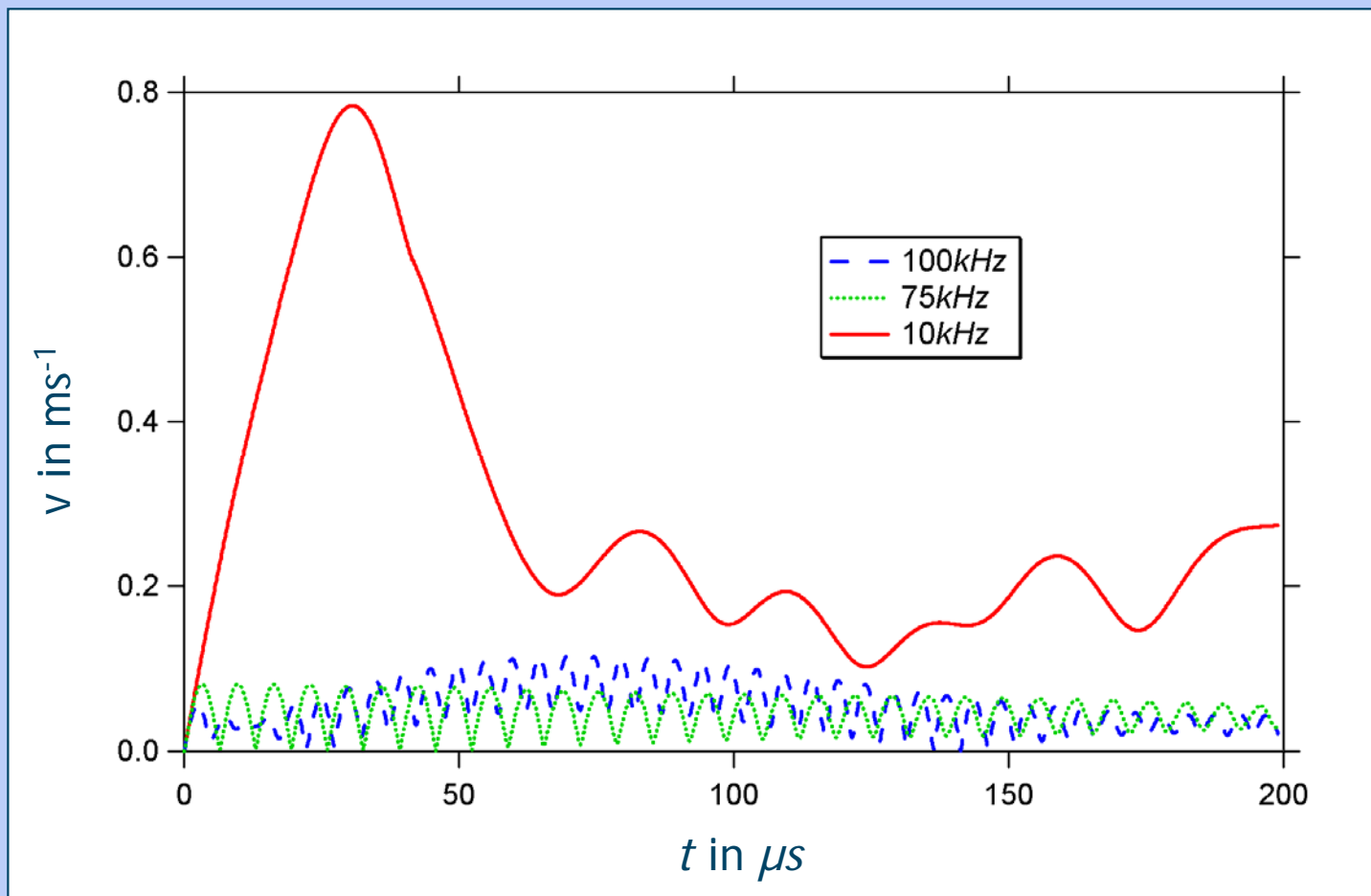
$$r = 20\mu\text{m}$$

Interactions of beads in fluids

15

Interactions of beads in fluids

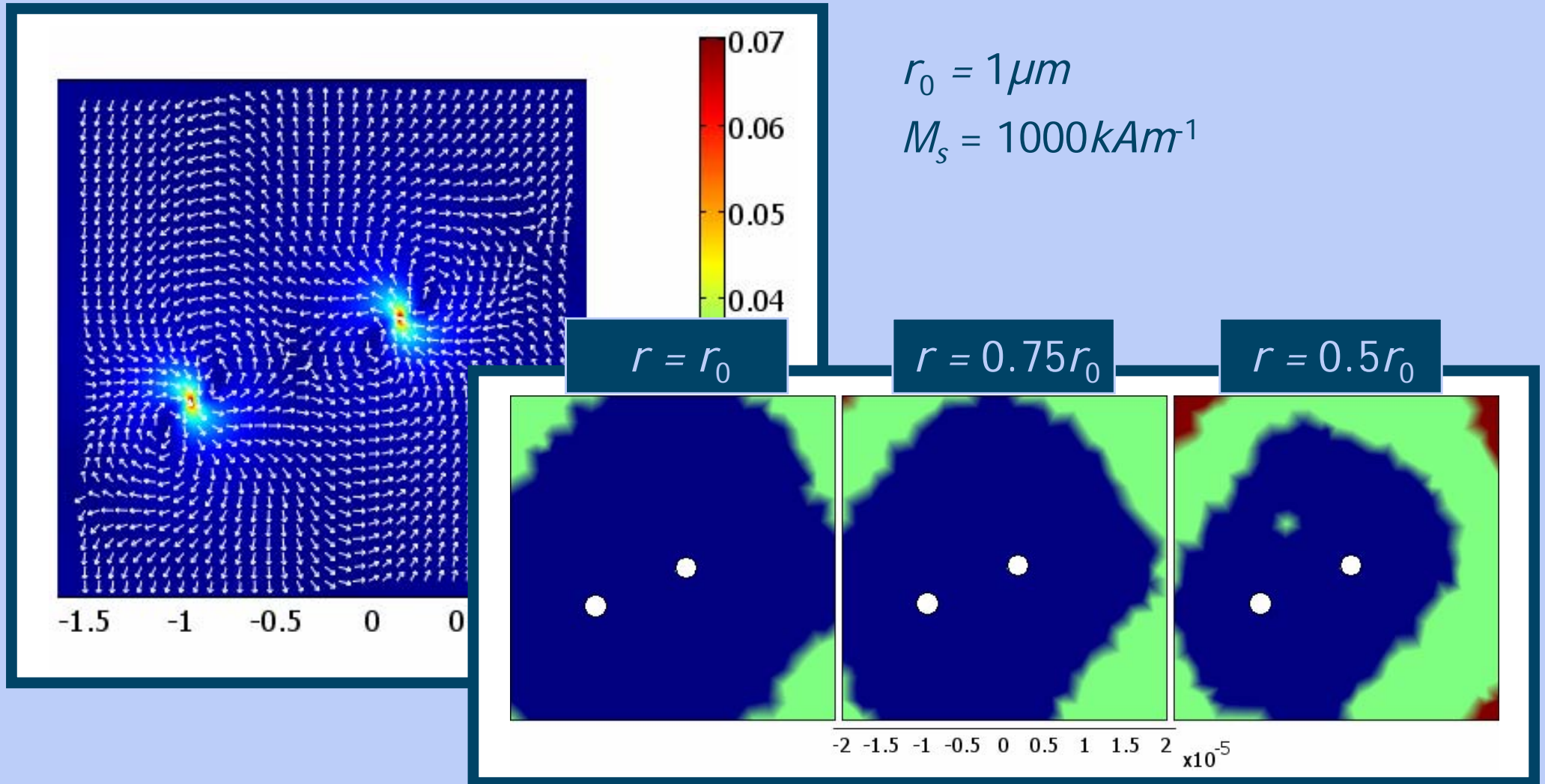
Observation:



$$M_s = 1000 \text{ kAm}^{-1}$$
$$r = 10 \mu\text{m}$$

high velocities might lead to non-laminar fluid behaviour on microscale!!

Comparison between forces



Conclusion & Outlook

Conclusion

- We have developed a model to describe the dynamic behaviour of magnetic beads
- We have simulated experimentally known effects (chain creation)
- We have shown that the magnetic interaction of particles can induce strong fluidic particle interactions that gain importance when dealing with different particle sizes

Outlook

- Finding proper clearcuts for different force regimes
- Implementing ferromagnetic particles