

COMPREHENSIVE NUMERICAL MODELING OF FILAMENTARY RRAM DEVICE.

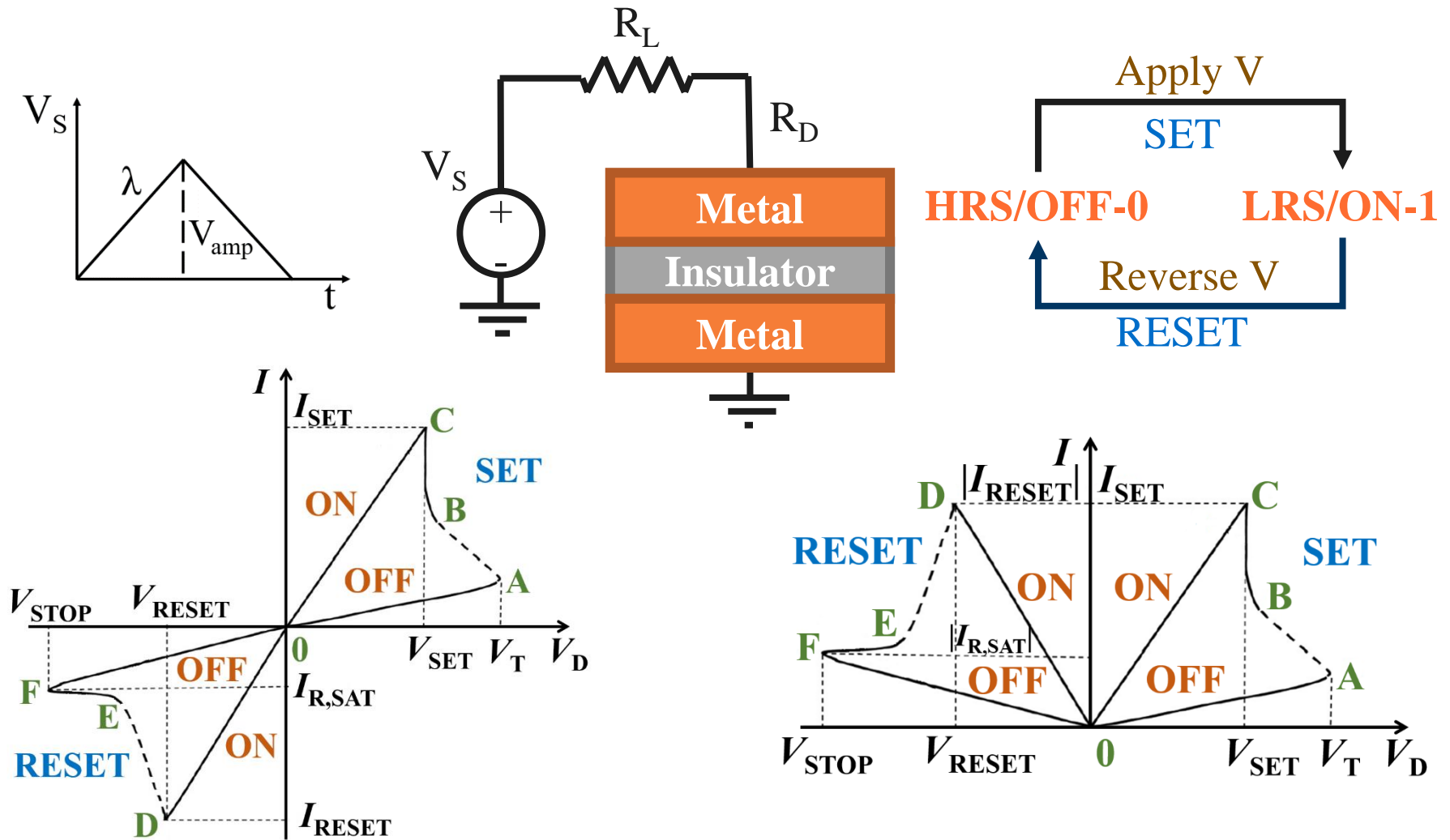
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COMSOL
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Research motivation: Bipolar resistive switching in RRAM



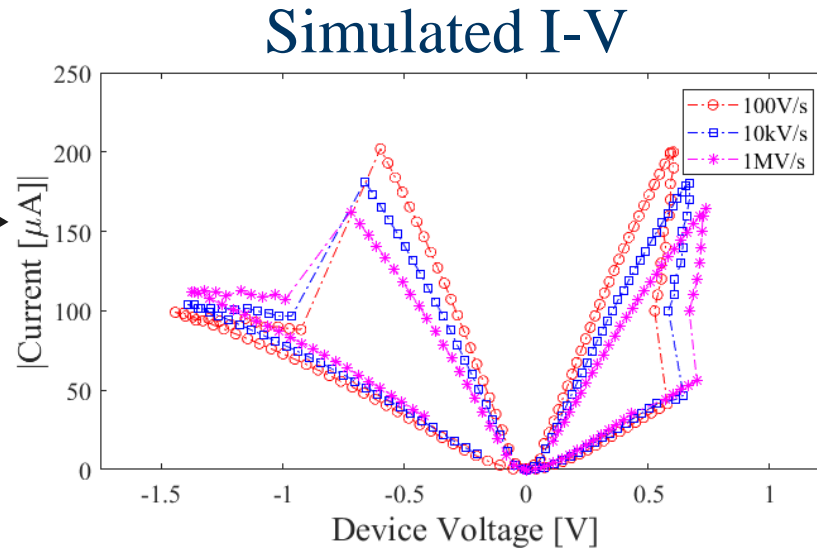
Statement of goal

- Develop a physics based numerical model of bipolar filamentary RRAM operation
 - independent of microscopic structure details
 - RRAM characteristics expressed through material parameters
 - I-V characteristics exhibiting *ramp-rate* and *cycle-to-cycle variations*

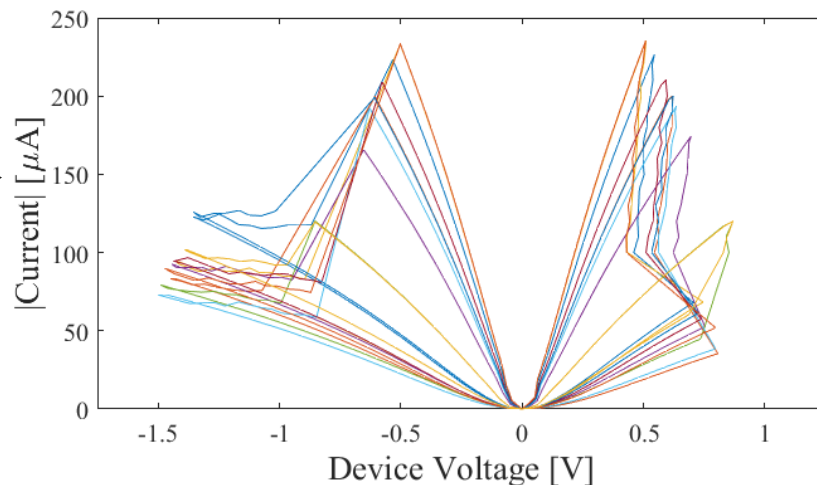
Showcasing our results

simulated I-V with ramp-rate and cycle-to-cycle variations

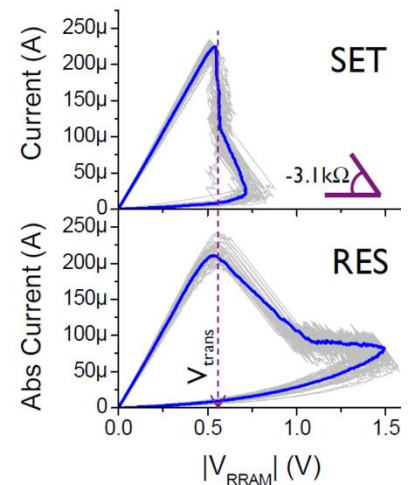
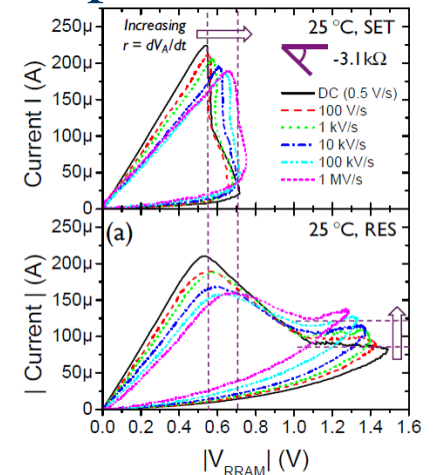
Increasing input voltage ramp-rate →



Multiple switching cycles →



Experimental I-V



Outline

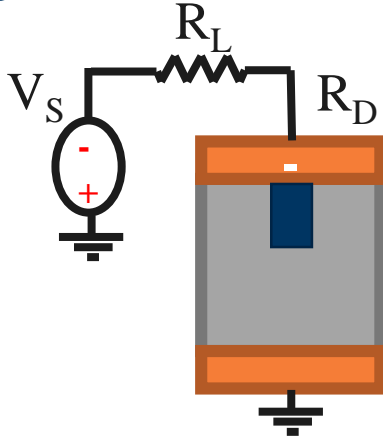
- Physics of Device operation
- Numerical Modeling
 - Program: Modules and Switching Condition
 - Device Model
 - Conduction Mechanism: Electric and Thermal
- Variability of RRAM parameters
 - Average ramp-rate dependent effects
 - Cycle-to-cycle Variations

Physics of Device operation:

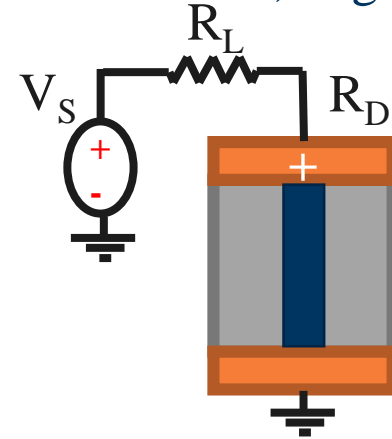
device operation consists of **Device States** and **Switching Processes**

Device States ('0' or '1')

OFF – Insulating State :
No or Partial Filament,
High Resistance, Low Current



ON – Conducting State:
Conducting Filament,
Low Resistance, High Current



Switching Processes (Write/Erase) :

SET : **OFF** → **ON**
Formation of
Conducting Filament

structural change

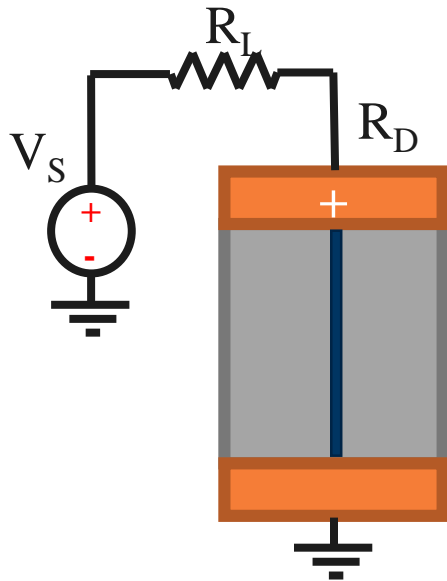
RESET : **ON** → **OFF**
Dissolution of
Conducting Filament

Physics Behind Switching Process:

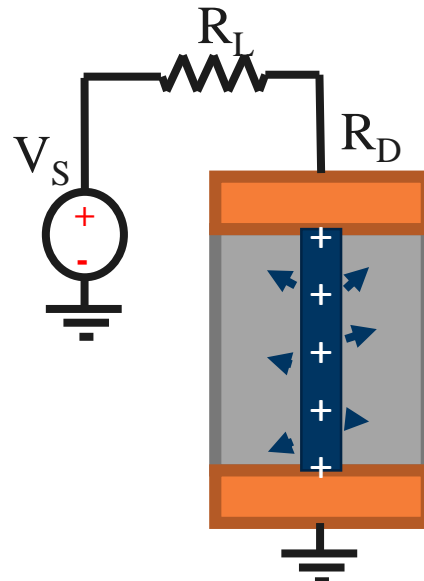
current carrying CF charges and produces strong radial field

$1\text{V}/10\text{nm} = 1\text{MV}/\text{cm}$ **SET**

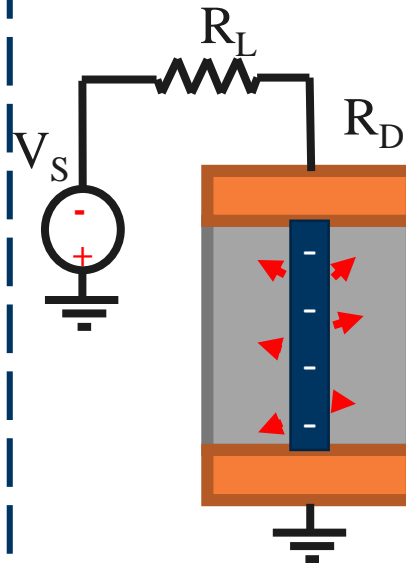
RESET



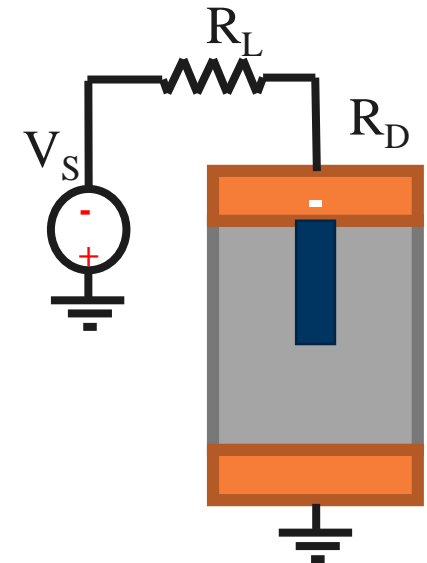
Field Induced
Nucleation then
shunting



CF charging
polarizes
insulating host
matrix



Reversing Polarity
charges CF
unfavorable to the
inherited polarization
of the host



Charged CF produces
a strong lateral field in
its vicinity opposite to
the host polarization,
then dissolves

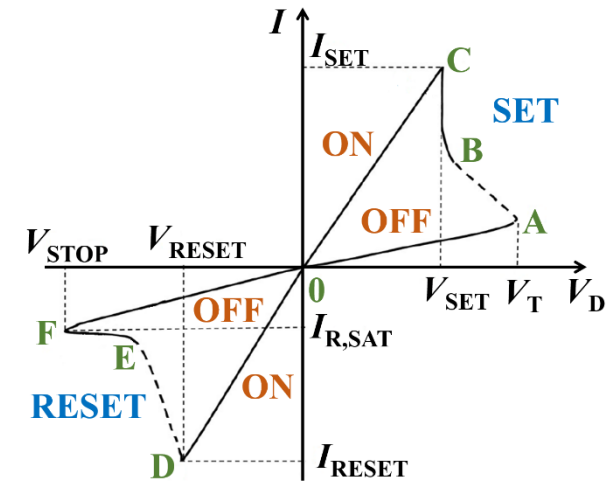
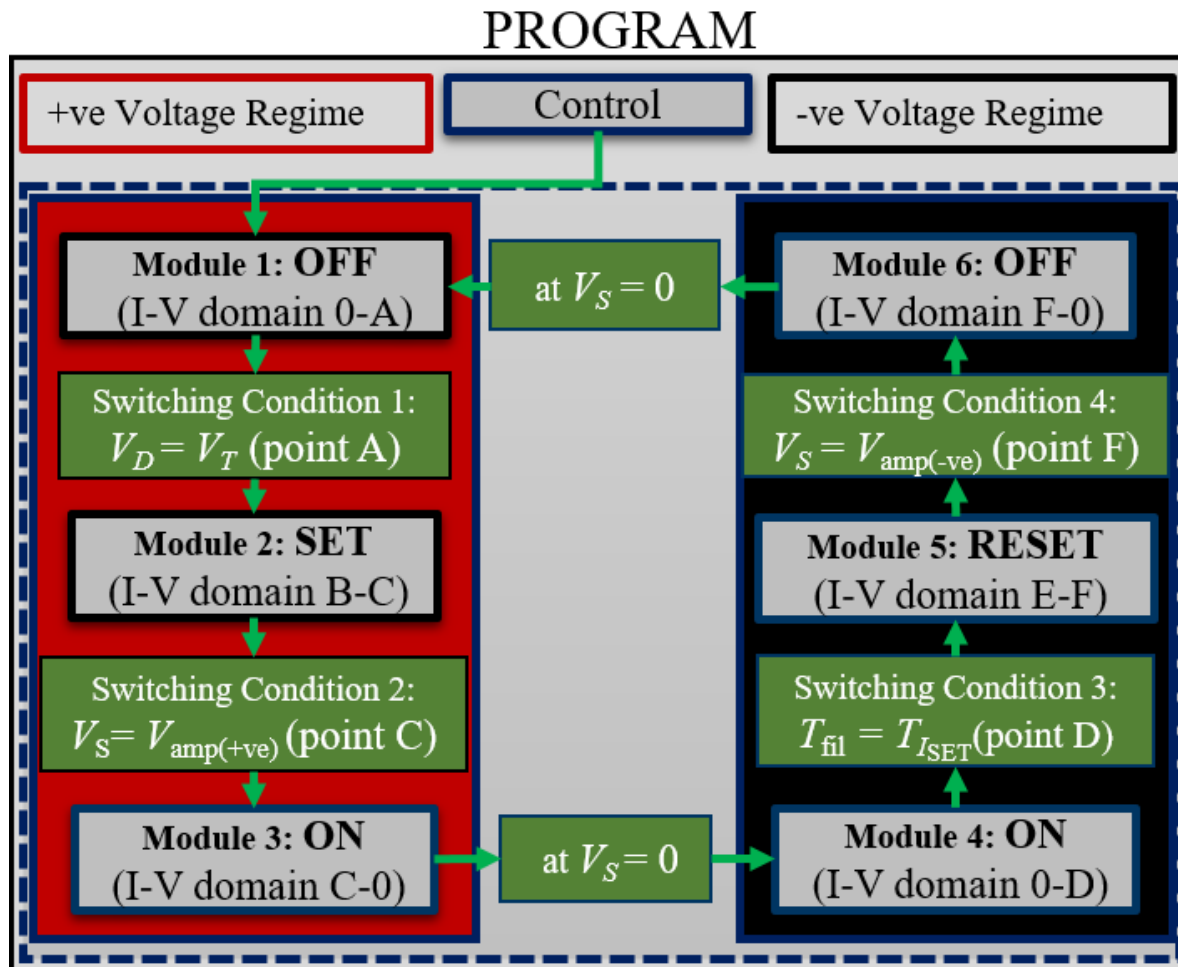
Note1: CF has finite capacitance

Note2: wire charging effect (due to Weber, 1852)

– overlooked in RRAM community

Numerical Modeling: Program

a MATLAB computer program is developed, consisting of 4 modules corresponding to the **device states** and **switching processes** executed sequentially mimicking actual device operation

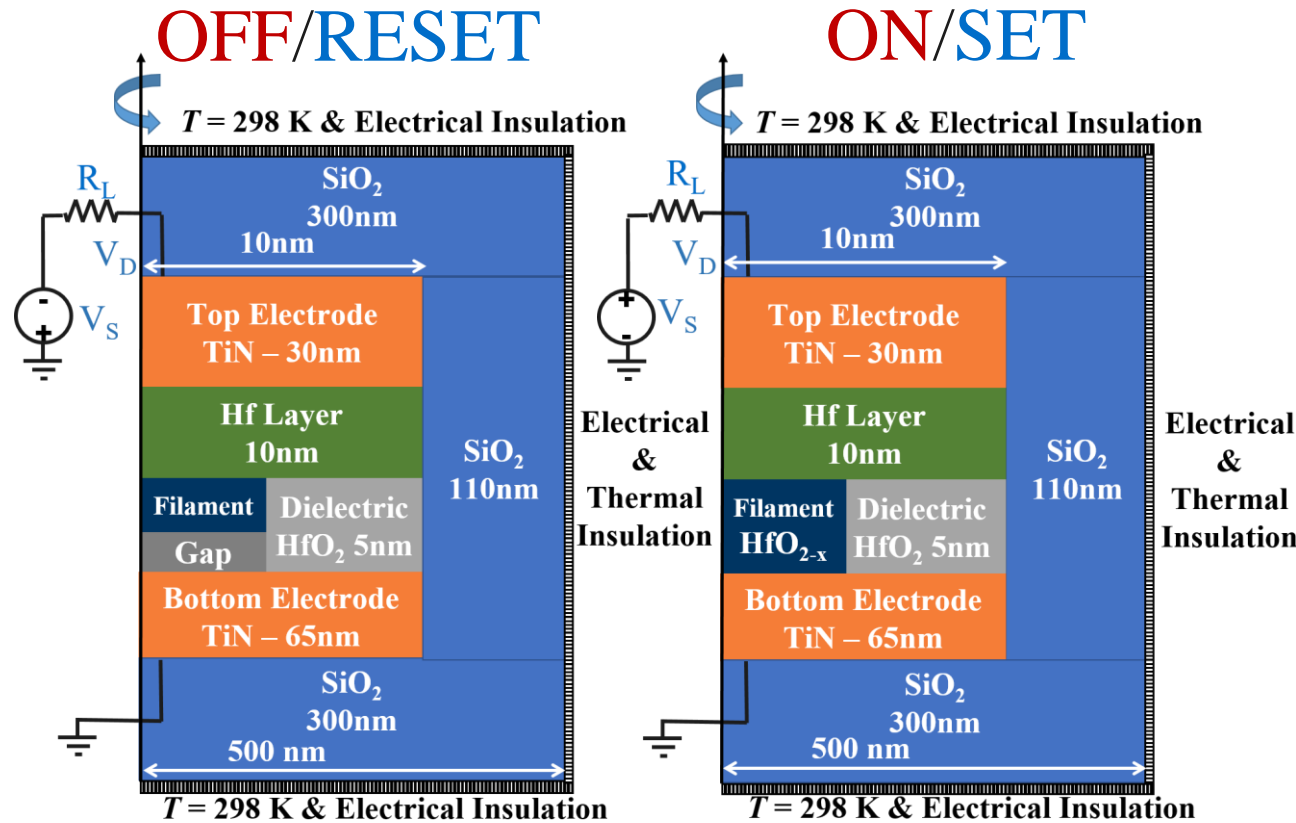


Device State: determined by conduction mechanism and filament dimensions

Switching process: governed by thermodynamics: (minimizing free energy by transforming phase)

Device Model:

TiN/Hf/HfO₂/TiN multilayered RRAM device is modeled in COMSOL



Utilized PDEs

OFF/ON	SET/RESET
<i>Electric Currents</i> module	
$\nabla \cdot \mathbf{J} = 0,$ $\mathbf{J} = \sigma_c \mathbf{E} + \epsilon \frac{\partial \mathbf{E}}{\partial t},$ $\mathbf{E} = -\nabla V$	$\nabla \cdot \mathbf{J} = 0,$ $\mathbf{J} = \sigma_c \mathbf{E},$ $\mathbf{E} = -\nabla V$
<i>Heat Transfer in Solids</i> module	
$\rho C_p \frac{\partial T}{\partial t}$ $-\nabla \cdot (\kappa \nabla T) = Q_s$	$-\nabla \cdot (\kappa \nabla T)$ $= Q_s$
<i>Multiphysics</i> module	
$Q_s = \mathbf{J} \cdot \mathbf{E}$	$Q_s = \mathbf{J} \cdot \mathbf{E}$
<i>Electric Circuit</i> module is used to define the circuitry	

OFF/ON: the gap/filament remains intact

RESET/SET: the gap/filament grows following the free energy minimum

Free Energy = Thermal + Electrostatic + Phase transition (Surface & Volume)

The thermal and electrostatic energies are obtained by solving coupled heat-electromagnetic PDE in COMSOL

Device Model:

material and various other parameters

Material	σ_c [S/m]	κ [W/K.m]	C_p [J/kg.K]	ϵ_r ^c	ρ [kg/m ³]	Parameter	Value	Parameter	Value
SiO ₂	10 ⁻⁹	1.38	703	3.9	2.2×10 ³	Circuitry		Chemical Energy	
TiN	Exp. $\sigma_c(T)$ ^a	$\sigma_c(T)TL$ ^d	545.33	-∞ ^f	5.22×10 ³	R_L	3.1 kΩ	σ	0.01 J/m ³
Hf	Exp. $\sigma_c(T)$ ^b	$\sigma_c(T)TL$ ^d	144	-∞ ^f	13.3×10 ³	$V_{amp(+ve)}$, $V_{amp(-ve)}$	1.25 V, -1.75 V	$\overline{\delta\mu_1}$	10 GJ/m ³
HfO ₂	10	0.5	120	25	10×10 ³	λ	100 V/s, 10 kV/s, 1 MV/s	$\overline{\delta\mu_2}$	6.5 GJ/m ³
HFO _{2-x}	$\sigma_{0f} \exp\left(-\alpha_f \ln\left(\frac{\tau}{\tau_0}\right)\right) \exp\left(\sqrt{\frac{eV}{kT}}\right)$	$\sigma_c(T)TL$ ^d	140 ^e	-∞ ^{e,f}	12×10 ^{3e}	Filament Nucleation		β_1	0.35 GJ/m ³
Gap	$\sigma_{0g} \exp\left(-\alpha_g \ln\left(\frac{\tau}{\tau_0}\right)\right) \exp\left(\sqrt{\frac{eV}{kT}}\right)$	$\kappa_{eff}\sigma_c(T)TL$ ^d	120	25	10×10 ³	h	5 nm	β_2	0.5 GJ/m ³
						W_0	2.5 eV	ΔW_{Buc}	1.0 eV
						Λ	6.6	ΔW_{Bi}	0.1 eV
						r_c	2.9 nm	ΔW_{Bmc}	0.3 eV
						r_{min}	0.5nm	Static Disorder	
						α	r_{min}/r_c	σ_{0f}	rand(2, 8) kS/m ^g
						Electric Conductivity		σ_{0g}	rand(1, 5) kS/m ^g
						σ_{0f}	5 kS/m	α_f	rand(-0.07, -0.03) ^g
						σ_{0g}	3 kS/m	α_g	rand(0.03, 0.07) ^g
						α_f	-0.05	$\overline{\delta\mu_1}$	rand(8.5, 11.5) GJ/m ^{3g}
						α_g	0.05	$\overline{\delta\mu_2}$	rand(5.5, 7.5) GJ/m ^{3g}
						τ	V_{amp}/λ	W_0	rand(2.4, 2.6) eV ^g
						$\tau_0(\tau_{min})$	0.1 ps	Thermal Conductivity	
								κ_{eff}	10
								^g function rand(x, y) produces uniformly distributed random number between x and y	

^a E. Langereis et al., *J. Appl. Phys.* **100**, 023534 (2006).

^b P. D. Desal, et al., *J. Phys. Chem. Ref. Data.* **3**, 1069 (1984).

^c Relative Permittivity

^d Wiedemann-Franz-Lorenz Law

^e Assumed value such that it lies in between Hf and HfO₂

^f -10⁶ was used instead of -∞ for practical purpose

M.A. Panzer, et al, *IEEE El. Dev. Lett.*, 30, pp. 1269-1271 (2009)

B. Govoreanu, et al., *IEEE Trans. El. Dev.*, 60, pp. 2471-2478 (2013)

E. Hildebrandt, et al., *Appl. Phys. Letts.*, 99, pp. 112902, (2011)

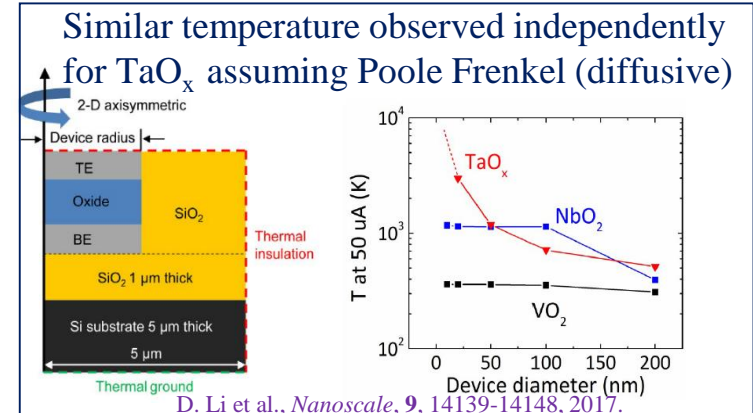
M.K. Samani, et al., *Thin Solids Films*, 573, pp. 108-112, (2013)

Carl L. Yaws. *The Yaws Handbook of Physical Properties for Hydrocarbons and Chemicals*, 2nd ed. (2015)

Conduction Mechanism:

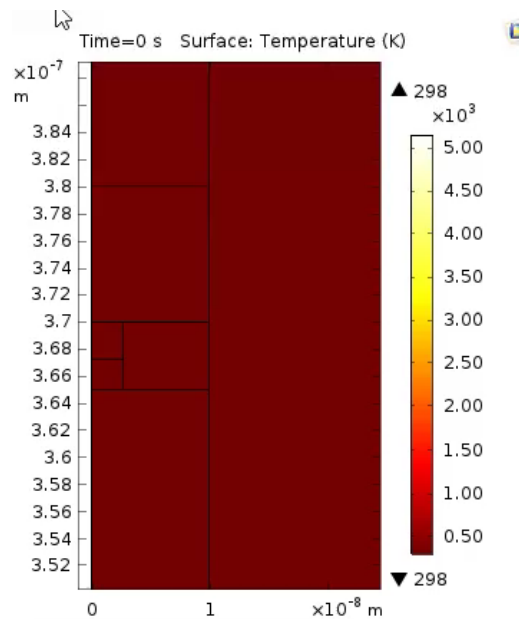
no consensus in RRAM community; we eliminate diffusion based mechanisms on the basis of unphysical device temperature

- Conduction mechanisms proposed: Poole-Frenkel, trap-assisted tunneling, Schottky emission, space charge limited current, variable range and nearest neighbor hopping
- All provide equally acceptable fits: non-indicative

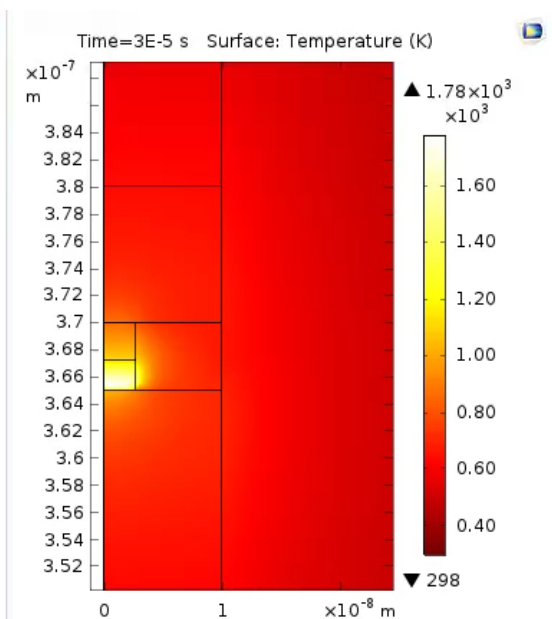


- Diffusion transport yields unphysical gap temperature ~ 4500K
- Electrons ballistically travels through ~2nm gap
- Correction term κ_{eff} introduced in thermal conductivity
- $\kappa_{\text{eff}} = \frac{\lambda_{\text{bal}}}{\lambda_{\text{diff}}} \approx \frac{10\text{nm}}{1\text{nm}}$
- $\kappa_{\text{gap}} = \kappa_{\text{eff}}\kappa_{\text{diff}}$
- λ is the mean free path

Diffusive heat Transport



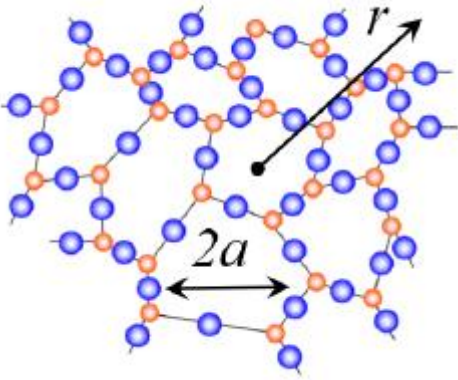
With Correction



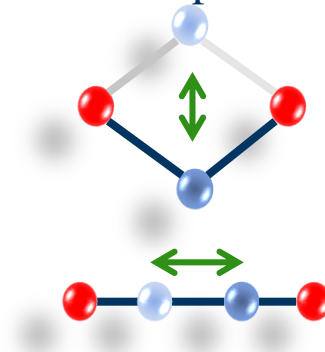
Variability of RRAM parameters:

HfO₂ is amorphous: some atoms retain mobility represented by random
Double Well Potentials (DWP)

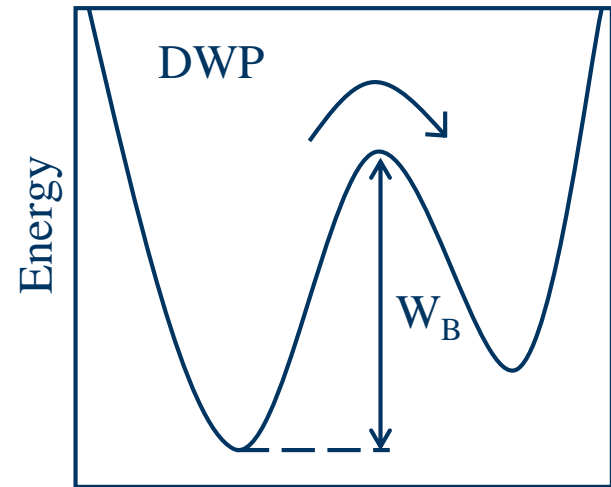
Disordered system: HfO₂



Blue Atom can reside
in either position



In noncrystalline structures some
atoms retain mobility moving in
random **Double Well Potentials**



Configurational coordinate

$$\text{Transition time } \tau = \exp\left(\frac{W_B}{kT}\right)$$

$$\text{Probability density } g(\tau) = \frac{1}{\tau}$$

$$\text{Cumulative density } \propto \ln \tau$$

Known since 1971 due to *Anderson, Halperin, Varma*
Verified via effects on:

Specific heat, Thermal expansion, Thermal
conductivity, Sound absorption, Sound propagation,
IR absorption, Light propagation, Point contact
resistance, Electronic noises, High pressure
behavior...more

Universal in noncrystalline metals, semiconductors,
and dielectrics

Average ramp-rate dependent effects

DWP affecting the electric conduction and energetics of filament and gap

The structural changes affects

- Electrical Conduction; the deformation due to DWP distorting the mobility edge and conduction;
- Chemical potential contribution

To Quantify

- the input voltage activates DWPs with transition time smaller than pulse time ($\tau < \tau_P$)
- Average fraction of activated DWP

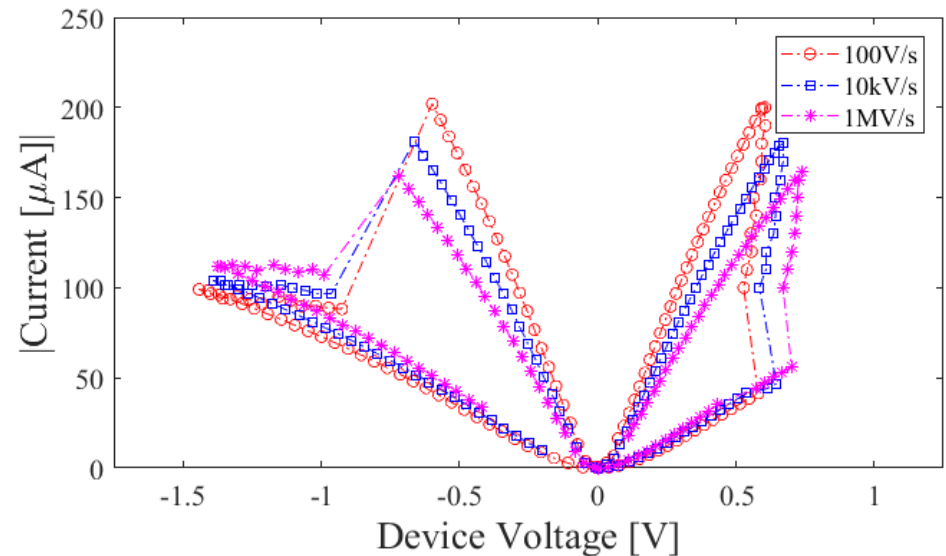
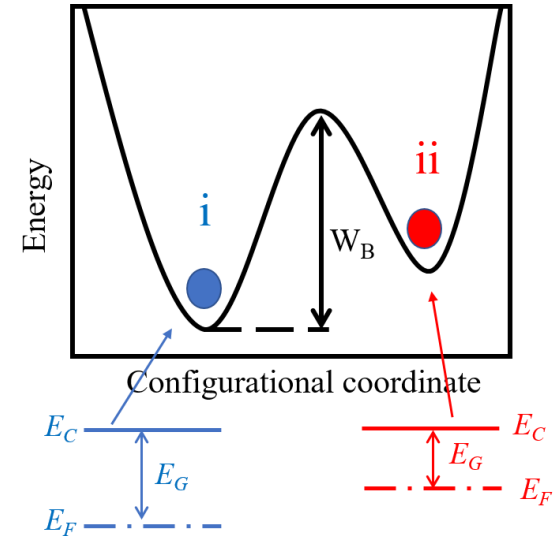
$$f(\tau_P) = \frac{kT}{\Delta W_B} \ln \left(\frac{\tau_P}{\tau_0} \right)$$

- Rate dependent Conduction

$$\sigma \sim \exp \left(\alpha \ln \left(\frac{\tau_P}{\tau_0} \right) \right), \text{ where } \alpha = \frac{u_0 D}{kT}$$

- Rate dependent Chemical Potential

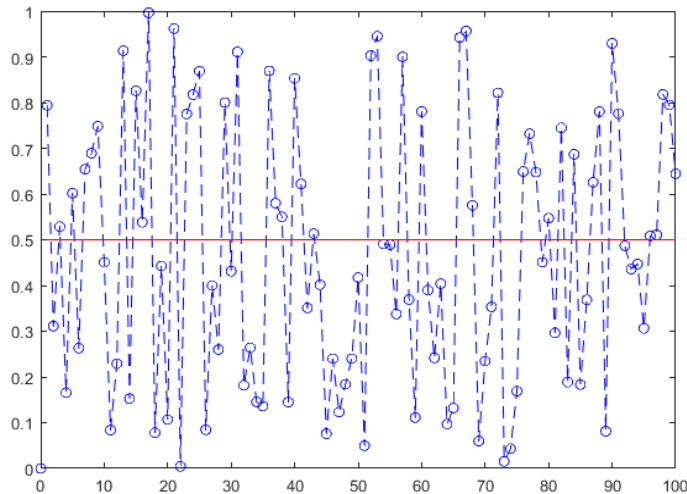
$$\delta\mu = \overline{\delta\mu} + \beta kT \left(\frac{1}{\Delta W_{uc}} - \frac{1}{\Delta W_i} \right) \ln \left(\frac{\tau_P}{\tau_0} \right)$$



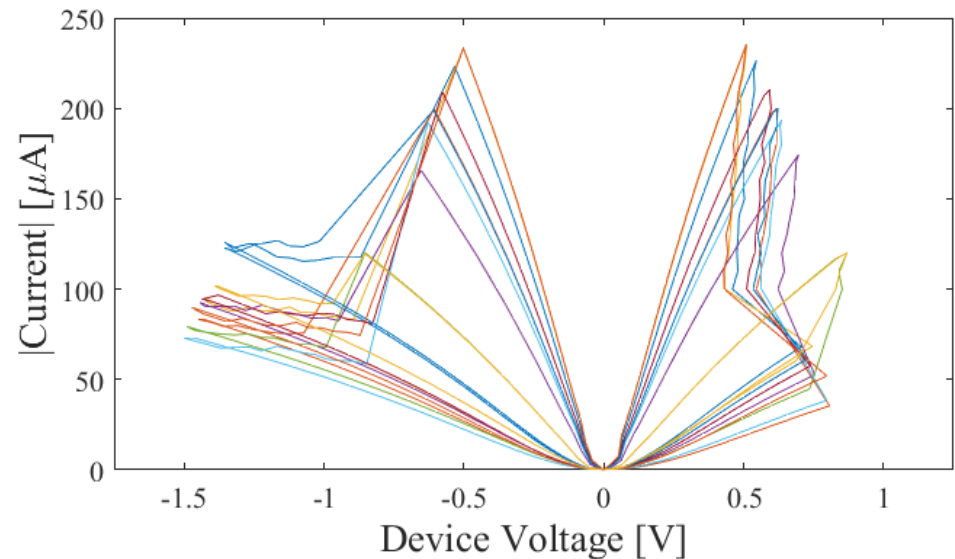
Cycle-to-cycle Variations

due to smallness of RRAM, the number of DWPs not enough to provide the statistical self-averaging

- Each switching cycle yields different structure
- The difference cause variation in DWP barrier heights, chemical potentials, and deformation potentials
- Applied random number generator to randomly vary DWP ensembles from once cycle to another



Uniformly distributed random number from 0 to 1



CONCLUSIONS

1. Developed physics based numerical model of RRAM including voltage ramp-rate and cycle-to-cycle variations
2. Established and incorporated the ballistic nature of electron transport
3. Included structural disorder of amorphous HfO_2 via double well atomic potentials modeling the observed variations between nominally identical RRAM devices
4. Developed a thermodynamic description of nano-sized modern RRAM and its numerical algorithm dramatically simpler than the kinetic approach

Acknowledgement

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- Collaborators
 - Ilya V. Karpov (Component Research, Intel)
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